

aminotetrazole (Britts and Karle¹⁷), the hydrazine salt of 5-aminotetrazole (Bryden¹⁸), and sodium tetrazolate monohydrate (Palenik¹⁹).

The nature of the Zn-Tz interaction may be summarized by comparing the results of this study with the results of the determination of the crystal structure of the iodopentamethylenetetrazolium monochloride complex (PMTICl) by Baenziger and Nelson.⁴

In both complexes the following statements apply. The moiety coordinated to the tetrazole ring is essen-

(17) K. Britts and I. L. Karle, *Acta Crystallogr.*, **22**, 308 (1967).

(18) J. H. Bryden, *ibid.*, **11**, 81 (1958).

(19) G. J. Palenik, *ibid.*, **16**, 596 (1963).

tially planar with the ring. The site of coordination to the tetrazole ring is identical (labeled N(4) in ZCTZ, N(2) in PMTICl). The coordination geometry at N(4) is essentially identical. Denoting values for the corresponding angles of PMTICl by an asterisk, the bond angles (degrees) about N(4) are as follows: N(3)-N(4)-C(5), 105 (2), 107 (2), 106* (4); N(3)-N(4)-Zn, 121 (2), 121 (2), 125* (4); C(5)-N(4)-Zn, 132 (2), 133 (2), 130* (4). These considerations suggest that the preferred mode of coordination in ZCTZ (and PMTICl) is by means of a charge-transfer σ bond to position N(4) of the tetrazole ring.

CONTRIBUTION FROM THE DEPARTMENT OF CHEMISTRY
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Optically Active Boron. The Crystal Structure of (+)-4-Methylpyridinetrimethylaminebromohydroboron Hexafluorophosphate

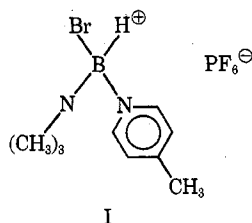
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The first crystal structure of a stable optically active compound having a boron atom as the asymmetric center has been elucidated. (+)-4-Methylpyridinetrimethylaminebromohydroboron hexafluorophosphate (I) crystallizes in the monoclinic system, space group $P2_1$, with $a = 19.35$ (1) Å, $b = 7.870$ (4) Å, $c = 10.549$ (5) Å, $\beta = 90^\circ 25'$ (10'). Two crystallographically independent formula units are present in the unit cell. The structure has been solved by Patterson and Fourier methods and refined to an R factor of 0.089 for 2481 independent nonzero reflections measured with counter technique (Cu $K\alpha$); the disagreement factor increases to 0.102 if 481 reflections below the limit of detectability are included. The geometrical parameters of the two independent cations, as well as of the two independent anions, are virtually identical. An almost regular sp^3 hybridization of the boron atom results from the arrangement of the surrounding ligands. The stereochemistry of the cations is discussed on the basis of intra- and intermolecular interactions; the presence of two pairs of independent ionic species is rationalized in terms of the tendency to produce a pseudocentrosymmetric packing.

Introduction

Recently Ryschkewitsch and Garrett have reported the preparation and the resolution of the first examples of stable optically active compounds with the boron atom as the asymmetric center.¹ While the isoelectronic and isosteric analogs of this compound, having optically active carbon or nitrogen atoms, have never been resolved, the boron salts have proven to be highly stable toward racemization during the fractionation process. In view of the interest involved in the stereochemical knowledge of such compounds, we have undertaken a single-crystal X-ray investigation of one of them, namely, (+)-4-methylpyridinetrimethylaminebromohydroboron hexafluorophosphate (I). We now



report the results of the detailed analysis of the crystal and molecular structure of I.

Experimental Section

Colorless, needlelike single crystals of I, obtained from an acetone-water mixture (mp 136-137°, $[\alpha]_D -60.9^\circ$), were kindly donated by Dr. Ryschkewitsch. A small crystal of approximately cylindrical shape ($r \approx 0.05$ mm, length ≈ 0.5 mm) was used for the X-ray investigation. From preliminary Weissenberg photographs the crystal was found to belong to monoclinic space group $P2_1$ (systematic $0k0$ absences with k odd). A Picker four-circle automated diffractometer equipped with a PDP-8 digital computer was used for the collection of the intensity data. The unit cell dimensions were obtained from a least-squares treatment of the 2θ , χ , and φ setting angles of 12 reflections with $2\theta > 80^\circ$, using a program prepared by Busing and Levy.² The resulting parameters are $a = 19.35 \pm 0.01$ Å, $b = 7.870 \pm 0.004$ Å, $c = 10.549 \pm 0.005$ Å, $\beta = 90^\circ 25' \pm 10'$, $d_{X\text{-ray}} = 1.607$ (1) g/cm³ with $Z = 4$, $d_{\text{exptl}} = 1.58$ (2) g/cm³ (floatation), $\mu = 51.56$ cm⁻¹, and $\lambda_{\text{Cu}K\alpha} 1.5418$ Å.

The experimental density agreed reasonably with the calculated density assuming four molecules of I per unit cell, thus requiring the presence of two crystallographically independent formula units. A total of 2962 integrated intensities were collected using a θ - 2θ scan mode (1.67°) with Ni-filtered Cu $K\alpha$ radiation in the range 0-130° of 2θ . Two stationary-crystal, stationary-counter background counts of 10 sec were taken at each end of each scan. Since $\mu = 51.66$ cm⁻¹ ($\lambda 1.5418$ Å) and

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(1) G. E. Ryschkewitsch and J. M. Garrett, *J. Amer. Chem. Soc.*, **90**, 7234 (1968).

(2) Z. W. R. Busing and H. A. Levy, *Acta Crystallogr.*, **22**, 457 (1967).

TABLE I: FINAL ATOMIC PARAMETERS^a

	A. Final Positional Parameters			A. Final Positional Parameters		
	1. Molecule a			2. Molecule b		
	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>
Br(1)	0.3937 (1)	0.0134 (4)	0.6857 (1)	0.0329 (1)	0.0000	-0.1441 (2)
P(1)	0.1769 (2)	0.2520 (6)	0.4518 (2)	0.3334 (2)	-0.3883 (5)	0.0547 (3)
F(1)	0.1145 (5)	0.2815 (19)	0.3610 (11)	0.4008 (5)	-0.4290 (19)	0.1364 (9)
F(2)	0.2046 (9)	0.1153 (26)	0.3581 (11)	0.3085 (6)	-0.2571 (17)	0.1597 (9)
F(3)	0.2149 (8)	0.3872 (25)	0.3715 (14)	0.2963 (7)	-0.5349 (17)	0.1273 (11)
F(4)	0.2429 (5)	0.2252 (16)	0.5414 (11)	0.2664 (5)	-0.3580 (21)	-0.0269 (9)
F(5)	0.1347 (7)	0.1217 (25)	0.5315 (11)	0.3712 (7)	-0.2462 (18)	-0.0184 (12)
F(6)	0.1528 (9)	0.3904 (25)	0.5416 (15)	0.3566 (6)	-0.5240 (16)	-0.0514 (18)
N(1)	0.4218 (5)	-0.0795 (15)	0.4201 (9)	0.0980 (5)	-0.0860 (15)	0.0978 (9)
N(2)	0.3476 (5)	-0.2995 (16)	0.5458 (9)	0.1662 (6)	0.1269 (16)	-0.0370 (10)
C(1)	0.4698 (7)	-0.1370 (19)	0.3406 (13)	0.1324 (6)	-0.2153 (19)	0.1554 (13)
C(2)	0.4793 (7)	-0.0693 (22)	0.2257 (12)	0.1175 (7)	-0.2667 (19)	0.2751 (13)
C(3)	0.4413 (7)	0.0701 (17)	0.1794 (11)	0.0669 (6)	-0.1922 (19)	0.3408 (11)
C(4)	0.3900 (22)	0.1234 (22)	0.2638 (12)	0.0273 (7)	-0.0615 (22)	0.2848 (14)
C(5)	0.3820 (6)	0.0553 (17)	0.3827 (11)	0.0442 (6)	-0.0171 (20)	0.1623 (12)
C(6)	0.4502 (10)	0.1396 (26)	0.0492 (15)	0.0479 (11)	-0.2439 (33)	0.4762 (17)
C(7)	0.3635 (9)	-0.4306 (23)	0.4485 (17)	0.2324 (8)	0.0945 (24)	0.0369 (17)
C(8)	0.2789 (8)	-0.2187 (28)	0.5092 (19)	0.1351 (11)	0.2840 (23)	0.0203 (18)
C(9)	0.3391 (11)	-0.3799 (30)	0.6743 (15)	0.1915 (10)	0.1658 (33)	-0.1728 (15)
B(1)	0.4098 (9)	-0.1686 (32)	0.5526 (14)	0.1209 (8)	-0.0409 (22)	-0.0401 (14)

B. Final Thermal Parameters $T = \exp[-1/4(B_{11}h^2a^{*2} + B_{22}k^2b^{*2} + B_{33}l^2c^{*2} + 2B_{12}hka^{*}b^{*} + 2B_{13}hla^{*}c^{*} + 2B_{23}k lb^{*}c^{*})]$

	1. Molecule a					
	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₁₂	<i>B</i> ₁₃	<i>B</i> ₂₃
Br(1)	10.9 (1)	9.5 (1)	4.5 (1)	-4.1 (1)	0.1 (1)	-1.3 (1)
P(1)	5.0 (1)	5.5 (2)	4.0 (1)	-0.2 (1)	0.1 (1)	0.0 (1)
F(1)	6.6 (5)	9.9 (7)	8.9 (6)	-0.1 (5)	-2.8 (4)	0.3 (6)
F(2)	17.4 (12)	14.2 (12)	5.9 (5)	5.6 (11)	-0.6 (6)	-2.2 (7)
F(3)	11.9 (9)	14.0 (12)	10.9 (8)	-4.0 (9)	-2.3 (7)	6.2 (9)
F(4)	5.7 (4)	8.7 (7)	8.9 (6)	-0.6 (5)	-1.8 (4)	0.3 (6)
F(5)	11.7 (8)	15.3 (13)	7.9 (6)	-6.1 (9)	-1.2 (6)	3.6 (8)
F(6)	15.4 (11)	13.0 (11)	11.5 (9)	7.0 (10)	-4.7 (8)	-5.4 (9)
N(1)	4.7 (4)	4.6 (5)	4.1 (4)	-0.6 (4)	-0.4 (3)	0.7 (4)
N(2)	4.0 (5)	6.1 (5)	3.3 (4)	-1.5 (4)	-0.7 (4)	1.0 (4)
C(1)	5.2 (6)	4.5 (6)	5.2 (6)	0.7 (5)	0.6 (4)	-1.2 (5)
C(2)	5.2 (6)	6.4 (8)	4.0 (5)	-0.2 (6)	0.2 (4)	-0.1 (6)
C(3)	5.5 (6)	4.5 (6)	3.7 (5)	-1.1 (5)	0.0 (4)	0.1 (4)
C(4)	4.9 (5)	6.2 (7)	4.0 (5)	-1.3 (6)	0.8 (4)	1.1 (5)
C(5)	4.6 (5)	4.8 (6)	4.0 (5)	-0.5 (5)	1.0 (4)	-0.6 (4)
C(6)	8.9 (10)	7.3 (10)	5.3 (7)	0.8 (8)	2.0 (7)	3.1 (7)
C(7)	7.5 (9)	5.8 (9)	7.5 (9)	-2.0 (7)	0.5 (7)	-1.2 (7)
C(8)	4.6 (6)	7.5 (10)	9.1 (11)	-1.7 (7)	-1.4 (6)	1.3 (9)
C(9)	10.0 (11)	8.5 (11)	4.7 (7)	-1.1 (10)	-0.2 (7)	2.3 (8)
B(1)	5.2 (7)	8.9 (12)	3.4 (6)	-0.4 (8)	-0.4 (5)	0.2 (7)

	2. Molecule b					
	<i>B</i> ₁₁	<i>B</i> ₂₂	<i>B</i> ₃₃	<i>B</i> ₁₂	<i>B</i> ₁₃	<i>B</i> ₂₃
Br(1)	6.7 (1)	10.9 (1)	6.1 (1)	-0.9 (1)	-2.3 (1)	1.8 (1)
P(1)	4.5 (1)	4.6 (1)	3.6 (1)	0.0 (1)	0.2 (1)	0.6 (1)
F(1)	7.0 (4)	12.5 (10)	6.5 (4)	4.0 (6)	-2.0 (4)	-0.4 (6)
F(2)	8.8 (5)	9.8 (7)	5.6 (4)	1.6 (6)	0.6 (4)	-1.7 (5)
F(3)	12.7 (8)	9.9 (9)	7.8 (6)	-6.3 (7)	0.1 (5)	2.5 (6)
F(4)	6.6 (4)	14.2 (10)	6.2 (4)	3.6 (6)	-2.7 (4)	-1.4 (6)
F(5)	11.7 (8)	8.9 (7)	9.2 (7)	-4.6 (7)	1.1 (6)	2.9 (6)
F(6)	11.2 (6)	8.6 (7)	5.0 (4)	4.2 (6)	0.6 (4)	-0.5 (4)
N(1)	3.7 (4)	4.5 (4)	4.1 (4)	0.0 (3)	-0.3 (3)	-0.3 (4)
N(2)	5.9 (5)	4.3 (5)	4.2 (4)	-0.7 (4)	0.5 (4)	0.5 (4)
C(1)	4.2 (5)	4.5 (6)	5.6 (6)	0.5 (5)	-0.8 (4)	0.1 (5)
C(2)	5.6 (6)	4.4 (6)	5.3 (6)	0.1 (5)	-0.7 (5)	0.4 (5)
C(3)	5.0 (5)	4.6 (6)	3.9 (5)	-1.3 (5)	1.0 (4)	-0.4 (4)
C(4)	5.4 (6)	6.2 (7)	5.1 (6)	0.6 (6)	0.4 (5)	-1.1 (6)
C(5)	4.4 (5)	4.9 (6)	5.2 (6)	0.7 (5)	0.5 (4)	-0.6 (5)
C(6)	9.0 (11)	9.8 (14)	5.7 (8)	-3.1 (11)	0.9 (7)	1.1 (9)
C(7)	5.2 (6)	6.2 (8)	7.5 (9)	0.7 (7)	-1.0 (6)	-0.2 (8)
C(8)	9.4 (11)	4.5 (7)	7.8 (9)	0.8 (8)	0.0 (8)	0.7 (7)
C(9)	9.6 (11)	11.6 (15)	4.1 (6)	-5.0 (11)	1.9 (7)	1.4 (8)
B(1)	4.6 (6)	5.0 (7)	4.4 (6)	0.4 (6)	0.1 (5)	0.0 (6)

^a Esd's in units of the last significant digit.

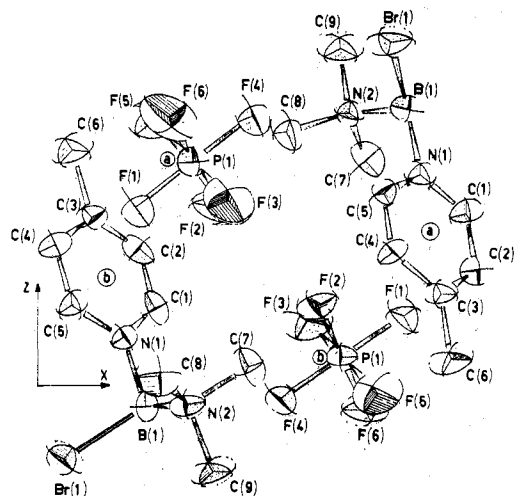


Figure 1.—ORTEP drawing of the thermal ellipsoids for the crystallographically independent unit, projected along the *b* axis. The two independent ions of each species are labeled as *a* and *b*.

proximately on the planes with $y = 0$ and $y = 1/2$; accordingly the positions of the two independent bromine atoms were located at nearly equal *y* coordinates thus giving rise to a quasicentrosymmetric arrangement. As a consequence the Fourier map calculated by phasing the reflections with the bromine atoms revealed the superimposition of two similar images of the structure, corresponding to a pseudosymmetry plane through the bromine atoms. However location of the two phosphorus atoms helped to remove part of the pseudosymmetry. The structure was then solved by three successive Fourier maps calculated after stepwise introduction of all the atoms. The refinement of the atomic parameters was carried out by means of seven cycles of 9×9 block-diagonal least-squares procedure,³ until the maximum shift in the atomic coordinates and anisotropic thermal factors was less than one-third of the corresponding standard deviation. The weighting scheme adopted throughout the refinement corresponds

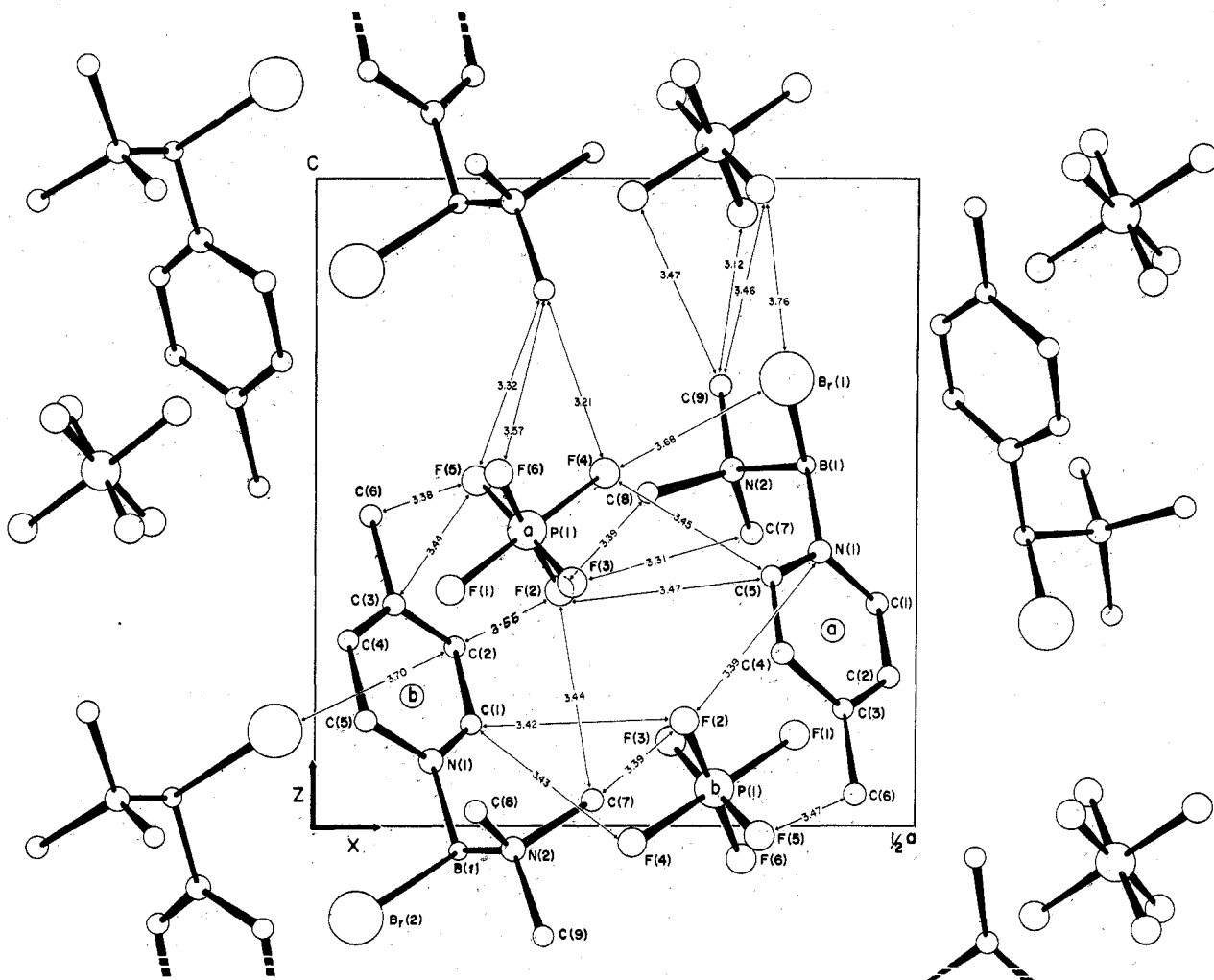


Figure 2.—Molecular packing of I, viewed along the *b* axis (cf. Figure 1). The shortest interatomic distances among different ions are indicated.

consequently $\mu R = 0.26$ (5), we felt it unnecessary to perform any absorption correction.

Structure Determination and Refinement

The structure was solved by the Patterson method. The heaviest interatomic vectors were seen to lie ap-

to that suggested by Cruickshank and Philling

$$w(hkl) = 1/a + bF_o(hkl) + cF_o^2(hkl)$$

(3) A complete set of crystallographic programs (FORTRAN IV) written by A. Immirzi was used throughout the determination and refinement of the structure: A. Immirzi, *Ric. Sci.*, **37**, 743, 846, 850 (1967).

TABLE II (Continued)

H	K	L	F0	FC	H	K	L	F0	FC	H	K	L	F0	FC	H	K	L	F0	FC	H	K	L	F0	FC
5	4	2	4	27	1	1	3	313	365	7	4	2	365	424	4	5	17	11	4	3	325	355		
5	4	3	4	27	0	1	3	485	491	7	4	3	202	222	8	4	6	220	192	9	1	2	295	295
5	4	4	4	27	0	1	3	580	585	7	4	4	128	118	8	4	6	286	276	9	1	2	295	295
5	4	5	4	27	0	1	3	680	685	7	4	5	86	86	8	4	6	336	317	9	1	2	295	295
5	4	6	4	27	0	1	3	780	785	7	4	6	44	44	8	4	6	386	367	9	1	2	295	295
5	4	7	4	27	0	1	3	880	885	7	4	7	22	22	8	4	6	436	417	9	1	2	295	295
5	4	8	4	27	0	1	3	980	985	7	4	8	10	10	8	4	6	486	467	9	1	2	295	295
5	4	9	4	27	0	1	3	1080	1085	7	4	9	8	8	8	4	6	536	517	9	1	2	295	295
5	4	10	4	27	0	1	3	1180	1185	7	4	10	6	6	6	4	6	586	567	9	1	2	295	295
5	4	11	4	27	0	1	3	1280	1285	7	4	11	4	4	4	4	6	636	617	9	1	2	295	295
5	4	12	4	27	0	1	3	1380	1385	7	4	12	2	2	2	2	6	686	667	9	1	2	295	295
5	4	13	4	27	0	1	3	1480	1485	7	4	13	1	1	1	1	6	736	717	9	1	2	295	295
5	4	14	4	27	0	1	3	1580	1585	7	4	14	0	0	0	0	6	786	767	9	1	2	295	295
5	4	15	4	27	0	1	3	1680	1685	7	4	15	0	0	0	0	6	836	817	9	1	2	295	295
5	4	16	4	27	0	1	3	1780	1785	7	4	16	0	0	0	0	6	886	867	9	1	2	295	295
5	4	17	4	27	0	1	3	1880	1885	7	4	17	0	0	0	0	6	936	917	9	1	2	295	295
5	4	18	4	27	0	1	3	1980	1985	7	4	18	0	0	0	0	6	986	967	9	1	2	295	295
5	4	19	4	27	0	1	3	2080	2085	7	4	19	0	0	0	0	6	1036	1017	9	1	2	295	295
5	4	20	4	27	0	1	3	2180	2185	7	4	20	0	0	0	0	6	1086	1067	9	1	2	295	295
5	4	21	4	27	0	1	3	2280	2285	7	4	21	0	0	0	0	6	1136	1117	9	1	2	295	295
5	4	22	4	27	0	1	3	2380	2385	7	4	22	0	0	0	0	6	1186	1167	9	1	2	295	295
5	4	23	4	27	0	1	3	2480	2485	7	4	23	0	0	0	0	6	1236	1217	9	1	2	295	295
5	4	24	4	27	0	1	3	2580	2585	7	4	24	0	0	0	0	6	1286	1267	9	1	2	295	295
5	4	25	4	27	0	1	3	2680	2685	7	4	25	0	0	0	0	6	1336	1317	9	1	2	295	295
5	4	26	4	27	0	1	3	2780	2785	7	4	26	0	0	0	0	6	1386	1367	9	1	2	295	295
5	4	27	4	27	0	1	3	2880	2885	7	4	27	0	0	0	0	6	1436	1417	9	1	2	295	295
5	4	28	4	27	0	1	3	2980	2985	7	4	28	0	0	0	0	6	1486	1467	9	1	2	295	295
5	4	29	4	27	0	1	3	3080	3085	7	4	29	0	0	0	0	6	1536	1517	9	1	2	295	295
5	4	30	4	27	0	1	3	3180	3185	7	4	30	0	0	0	0	6	1586	1567	9	1	2	295	295
5	4	31	4	27	0	1	3	3280	3285	7	4	31	0	0	0	0	6	1636	1617	9	1	2	295	295
5	4	32	4	27	0	1	3	3380	3385	7	4	32	0	0	0	0	6	1686	1667	9	1	2	295	295
5	4	33	4	27	0	1	3	3480	3485	7	4	33	0	0	0	0	6	1736	1717	9	1	2	295	295
5	4	34	4	27	0	1	3	3580	3585	7	4	34	0	0	0	0	6	1786	1767	9	1	2	295	295
5	4	35	4	27	0	1	3	3680	3685	7	4	35	0	0	0	0	6	1836	1817	9	1	2	295	295
5	4	36	4	27	0	1	3	3780	3785	7	4	36	0	0	0	0	6	1886	1867	9	1	2	295	295
5	4	37	4	27	0	1	3	3880	3885	7	4	37	0	0	0	0	6	1936	1917	9	1	2	295	295
5	4	38	4	27	0	1	3	3980	3985	7	4	38	0	0	0	0	6	1986	1967	9	1	2	295	295
5	4	39	4	27	0	1	3	4080	4085	7	4	39	0	0	0	0	6	2036	2017	9	1	2	295	295
5	4	40	4	27	0	1	3	4180	4185	7	4	40	0	0	0	0	6	2086	2067	9	1	2	295	295
5	4	41	4	27	0	1	3	4280	4285	7	4	41	0	0	0	0	6	2136	2117	9	1	2	295	295
5	4	42	4	27	0	1	3	4380	4385	7	4	42	0	0	0	0	6	2186	2167	9	1	2	295	295
5	4	43	4	27	0	1	3	4480	4485	7	4	43	0	0	0	0	6	2236	2217	9	1	2	295	295
5	4	44	4	27	0	1	3	4580	4585	7	4	44	0	0	0	0	6	2286	2267	9	1	2	295	295
5	4	45	4	27	0	1	3	4680	4685	7	4	45	0	0	0	0	6	2336	2317	9	1	2	295	295
5	4	46	4	27	0	1	3	4780	4785	7	4	46	0	0	0	0	6	2386	2367	9	1	2	295	295
5	4	47	4	27	0	1	3	4880	4885	7	4	47	0	0	0	0	6	2436	2417	9	1	2	295	295
5	4	48	4	27	0	1	3	4980	4985	7	4	48	0	0	0	0	6	2486	2467	9	1	2	295	295
5	4	49	4	27	0	1	3	5080	5085	7	4	49	0	0	0	0	6	2536	2517	9	1	2	295	295
5	4	50	4	27	0	1	3	5180	5185	7	4	50	0	0	0	0	6	2586	2567	9	1	2	295	295

in Table I together with their corresponding standard deviations; Figure 1 shows an ORTEP drawing of the thermal ellipsoids. A list of the observed and calculated structure factors is also given in Table II, while Table III lists bond lengths, bond angles, and some relevant dihedral angles for the independent unit of the structure.

Discussion of the Structure

The presence of two chemically identical but crystallographically independent formula units in the unit cell, a and b in Figures 1 and 2, is rather unusual. However, as can be seen from the figures, a pseudocentrosymmetric arrangement of the ions exists around the points with approximate fractional coordinates (1/4, -0.06, 1/4), (1/4, -0.06, 3/4), and the symmetry-related ones. While the PF₆⁻ ions are in good agreement with such pseudosymmetry, the cations depart from it as far as the positions in space of the bromine atoms are concerned. This departure results from the presence of optically active cations having all the same chirality, as it is unambiguously derived from the spatial distribution of the three Fourier-detectable substituents attached to the boron atoms (the hydrogen atom must necessarily occupy the remaining tetrahedral site). A virtually complete centrosymmetry could be achieved only through the exchange of Br with H in the coordination sites around the boron atom in either of the two independent cations, producing a racemic mixture. We believe that the present structure represents an interesting example in which the tendency toward the retention of a center of symmetry—which in most cases is a favor-

able condition of good packing—associated with the presence of chiral isomorphous molecules, is the driving force toward the occurrence of two independent molecules in the asymmetric unit.

The geometrical parameters (bond lengths, bond angles, and dihedral angles) of the four independent ionic species are reported in Table III. Neither of the two pairs of independent ions shows any substantial differences in geometry: in Figure 3 the average dimen-

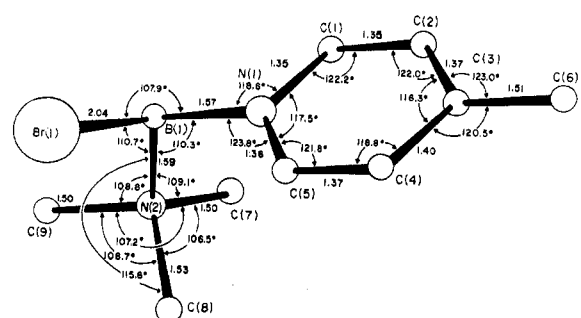


Figure 3.—Molecular geometry of the cation; the parameters are averaged over the two crystallographically independent individuals. For the esd and the dihedral angles, cf. Table III.

sions of the cation are reported. As for the hexafluorophosphate anion, its geometry can be described as a slightly distorted octahedron (cf. Table III), with a d_{P-F} average value of 1.57 ± 0.04 Å.⁵ The two chemically different B-N bond lengths do not show apprecia-

(5) In the present paper the deviations associated with any average value indicate the range of the individual values.

TABLE III
MOLECULAR DIMENSIONS^a

	Molecule a	Molecule b
A. Bond Lengths, Å		
Br(1)-B(1)	2.03 (2)	2.05 (2)
P(1)-F(1)	1.55 (1)	1.59 (1)
P(1)-F(2)	1.56 (2)	1.59 (1)
P(1)-F(3)	1.55 (2)	1.56 (1)
P(1)-F(4)	1.60 (1)	1.57 (1)
P(1)-F(5)	1.56 (2)	1.55 (1)
P(1)-F(6)	1.52 (2)	1.61 (1)
N(1)-B(1)	1.58 (2)	1.56 (2)
N(1)-C(1)	1.34 (2)	1.36 (2)
N(1)-C(5)	1.37 (2)	1.36 (2)
N(2)-B(1)	1.59 (2)	1.59 (2)
N(2)-C(7)	1.49 (2)	1.52 (2)
N(2)-C(8)	1.52 (2)	1.55 (2)
N(2)-C(9)	1.51 (2)	1.50 (2)
C(1)-C(2)	1.34 (2)	1.36 (2)
C(2)-C(3)	1.41 (2)	1.34 (2)
C(3)-C(4)	1.40 (2)	1.41 (2)
C(4)-C(5)	1.37 (2)	1.38 (2)
C(3)-C(6)	1.49 (2)	1.53 (2)
B. Bond Angles, Deg		
F(1)-P(1)-F(2)	88.9 (4)	90.3 (3)
F(1)-P(1)-F(3)	86.0 (4)	87.9 (3)
F(1)-P(1)-F(4)	177.8 (4)	177.1 (4)
F(1)-P(1)-F(5)	91.3 (4)	91.6 (3)
F(1)-P(1)-F(6)	92.1 (4)	90.7 (3)
F(2)-P(1)-F(3)	87.7 (5)	89.7 (3)
F(2)-P(1)-F(4)	90.4 (3)	91.7 (3)
F(2)-P(1)-F(5)	94.1 (4)	91.4 (4)
F(2)-P(1)-F(6)	177.3 (4)	178.4 (4)
F(3)-P(1)-F(4)	92.0 (4)	90.0 (3)
F(3)-P(1)-F(5)	176.6 (4)	178.8 (4)
F(3)-P(1)-F(6)	89.8 (5)	89.0 (3)
F(4)-P(1)-F(5)	90.8 (4)	90.5 (4)
F(4)-P(1)-F(6)	88.5 (5)	87.3 (3)
F(5)-P(1)-F(6)	88.3 (5)	89.9 (3)
Br(1)-B(1)-N(1)	108.8 (3)	107.1 (3)
Br(1)-B(1)-N(2)	111.7 (6)	109.7 (5)
N(1)-B(1)-N(2)	111.4 (5)	109.3 (5)
B(1)-N(2)-C(7)	108.7 (6)	109.5 (6)
B(1)-N(2)-C(8)	113.7 (6)	118.0 (7)
B(1)-N(2)-C(9)	108.7 (7)	109.0 (6)
B(1)-N(1)-C(1)	120.8 (5)	116.4 (5)
B(1)-N(1)-C(5)	120.9 (6)	126.7 (6)
C(7)-N(2)-C(8)	107.4 (6)	105.7 (6)
C(7)-N(2)-C(9)	110.7 (6)	103.7 (5)
C(8)-N(2)-C(9)	107.6 (6)	109.9 (6)
N(1)-C(1)-C(2)	122.3 (7)	122.2 (6)
C(1)-C(2)-C(3)	123.4 (6)	120.7 (6)
C(2)-C(3)-C(4)	112.6 (6)	120.1 (7)
C(3)-C(4)-C(5)	123.2 (6)	116.4 (6)
C(4)-C(5)-N(1)	120.0 (5)	123.7 (6)
C(5)-N(1)-C(1)	118.3 (6)	116.7 (6)
C(2)-C(3)-C(6)	122.9 (7)	123.2 (7)
C(4)-C(3)-C(6)	124.3 (7)	116.7 (7)
C. Internal Rotation Angles, Deg (Trans Conformation 180°)		
N(1)-B(1)-N(2)-C(7)	59	65
N(1)-B(1)-N(2)-C(8)	-60	-56
N(1)-B(1)-N(2)-C(9)	180	178
Br(1)-B(1)-N(2)-C(7)	179	178
Br(1)-B(1)-N(2)-C(8)	61	61
Br(1)-B(1)-N(2)-C(9)	-58	-66
C(1)-N(1)-B(1)-Br(1)	140	138
C(5)-N(1)-B(1)-Br(1)	-42	-38
C(1)-N(1)-B(1)-N(2)	-97	-104
C(5)-N(1)-B(1)-N(2)	81	81

^a Esd's in units of the last significant digit.

ble difference, and their average value (1.579 ± 0.015 Å) is in the range of values quoted in the literature for B-N bond lengths involving tetrahedral boron atoms.⁶ To the best of our knowledge, the B-Br distance (average 2.038 ± 0.008 Å) is the first reported bond length of this type involving a B atom with sp^3 hybridization, and it is appreciably longer than the B-Br distance reported for the gaseous BBr_3 molecule (1.87 ± 0.02 Å).⁷ The N-B-N and N-B-Br bond angles are all close to the tetrahedral value (within $\pm 2.5^\circ$), in agreement with the sp^3 hybridization on the boron atom. The systematic increase of $1-2^\circ$ of these bond angles in ion b compared to a is probably significant (the average esd on the angles is 0.4°) and is to be attributed to packing effects. In both cations, the observed compression of the Br-B-N (picolinic) angle is apparently due to the steric interactions involving the bulky $-N(CH_3)_3$ group. One would expect that the nonbonded interactions around the boron atom would give rise to some increase of the N-B-N and N-B-Br angles with respect to the tetrahedral value. The observed values, however, do not differ on the average. This may be rationalized as follows. Since B-H is certainly the shortest of the four B-X bonds, the repulsion between the bonding electron pairs around the boron atom is responsible for some compression effect, which leads to the observed geometry (Sidgwick and Powell's rule).⁸ The rotation around the B-N(CH_3)₃ bond (*cf.* Table III and Figure 4) obtains a nearly staggered arrangement of the sub-

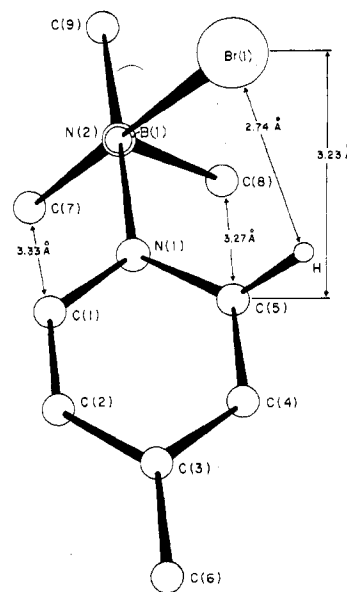


Figure 4.—View of the cation (average geometry) along the B(1)-N(2) bond. The shortest intramolecular distances between atoms separated by more than three bonds are indicated. Only one hydrogen atom has been represented (see text).

stituents around the B and N atoms. The CH_3-N-CH_3 bond angles are slightly smaller than the tetrahedral

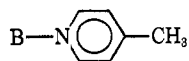
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value ($107.5 \pm 3.5^\circ$), only one of the six independent measurements being larger than 109.5° (see Table III). The C(8)-N-B angle (average $116.0 \pm 2.1^\circ$) is significantly larger than the other two corresponding angles (average $108.9 \pm 0.8^\circ$), apparently because of the larger steric repulsions involving the C(8) group, which is in "gauche" orientation with respect to both of the bulky substituents on the boron atom.

The



group is planar within experimental error. The msd's from the average plane are 0.018 and 0.027 Å for cations a and b, respectively; the equations of the corresponding planes, in terms of fractional coordinates, are $0.6388x + 0.6753y + 0.3640z - 1.1072 = 0$, for a, and $0.6611x + 0.6389y + 0.3885z - 6.7044 = 0$, for b. The dihedral angles between the picolinic group and the Br-B-N(1) and N(2)-B-N(1) planes are (average) 40.0 ± 2.5 and $81.0 \pm 0.3^\circ$, respectively; the larger dihedral angle involving the bulkier $-N(CH_3)_2$ group is very likely due to larger steric repulsions. As it is shown in Figure 4, the pyridine ring is nearly orthogonal to the N-B-N plane (average value of the dihedral angle $82 \pm 1^\circ$); the departure from the 90° value seems to be due to the necessity of increasing the intramolecular Br...H dis-

tance (2.7 Å, see Figure 4; the distance has been evaluated assuming the H atom in the normal stereochemical position), which still falls short of the sum of the van der Waals radii (3.1 Å).⁹

The geometry of the pyridine ring of the picolinic residue does not differ, within experimental error, from that established for pyridine by microwave measurements.¹⁰ As expected in an ionic structure, ions with the same charge tend to avoid close contacts. As reported in Figure 2, where hydrogen atoms are neglected, the shortest interatomic distance between different cations is 3.70 Å (Br...C); as for the PF_6^- ions the F...F distances are never shorter than 4 Å. On the other hand, interatomic contact between ions of opposite charge can be as close as 3.12 Å (F...CH₃) and 3.39 Å (F...N of the picolinic ring).

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The Crystal and Molecular Structure of Dimeric Tetramethyl-1,3-diisothiocyanatodistannoxane, $[((CH_3)_2SnNCS)_2O]_2$

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The crystal structure of dimeric tetramethyl-1,3-diisothiocyanatodistannoxane has been determined from three-dimensional X-ray diffraction data. It forms monoclinic crystals in space group $C2/m$ with $a = 21.73$ (4) Å, $b = 7.88$ (1) Å, $c = 10.96$ (2) Å, and $\beta = 134.7$ (2) $^\circ$. The intensities of 546 independent reflections were obtained by visual estimation of film data. The structure was solved by Patterson and Fourier syntheses and refined by the least-squares method to a conventional R factor of 8.0%. The dimeric molecule has $2/m$ symmetry with a central planar four-membered $Sn(2)_2O_2$ ring. Secondary bonding (2.84 (4) Å) between the isothiocyanate nitrogens of these tin atoms and the exocyclic tins forms two additional planar $Sn(2)OSn(1)\cdots N(2)$ rings. The endocyclic tin is thus pentacoordinate, approximately trigonal bipyramidal. Still further weak $Sn(1)\cdots S(2)$ linkage (3.33 (3) Å) between dimers results in distorted octahedral hexacoordination for the exocyclic tin atom and infinite chains parallel to [001]. Distances are $Sn(1)-O = 2.05$ (5), $Sn(1)-N(1) = 2.06$ (4), $Sn(1)-CH_3(1) = 2.07$ (3), $Sn(2)-O = 1.99$ (3), $Sn(2)-O = 2.15$ (4), $Sn(2)-N(2) = 2.39$ (8), and $Sn(2)-CH_3(2) = 2.15$ (3) Å. Both NCS groups are linear with distances $N(1)-C(1) = 1.17$ (7), $C(1)-S(1) = 1.61$ (6), $N(2)-C(2) = 1.10$ (7), and $C(2)-S(2) = 1.64$ (5) Å.

Introduction

The crystal structure reported here is a part of a general investigation of donor-acceptor interactions between pseudohalogens and heavy atoms in the solid state. Many tetraalkyldistannoxane derivatives $(XR_2-Sn)_2O$ and the related hydroxides $XR_2SnOSnR_2OH$ are known ($X =$ halogen, pseudohalogen, carboxylate, ni-

trate, alkoxide, or phenoxide). Pfeiffer and Brack¹ in 1914 assumed that the compound $(XR_2Sn)_2O$ had the structure I; however, Harada² argued the case for a cyclic structure II where the tin atoms have coordination numbers of 4 and 5; on rather insecure evidence he

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